



COHERENCE QED IN COLD FUSION

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Introduction

In the usual QED of Condensed Matter the start point is:

$$H_{\text{matt}} = \sum_{j=1}^N H_0(\vec{x}_j, \vec{p}_j, \alpha_j) + V(\vec{x}_1, \alpha_1, \dots, \vec{x}_N, \alpha_N)$$

where (generally) V is a short range potential.

Introducing the orthonormal complete set of vector that diagonalizes H_0 ,

we have

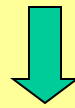
$$\begin{array}{ccc} \{\phi_n(\vec{x}, \alpha)\} & \xrightarrow{\text{Fock-space}} & H_0 = \sum_{j=1}^N H_0(\vec{x}_j, \vec{p}_j, \alpha_j) = \sum_n \varepsilon_n a_n^+ a_n \\ \downarrow & & \\ E\{n_m\} = \sum_m n_m \varepsilon_m & & \Psi(\vec{x}, \alpha, t) = \sum_n a_n(t) \phi_n(\vec{x}, \alpha) \end{array}$$

QED Lagrangian

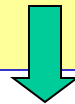
By means of field operator we have

$$L_{\text{matt}} = i \int_{\vec{x}, \alpha} \Psi^+(\vec{x}, \alpha, t) \Psi_t(\vec{x}, \alpha, t) - H_0 \quad [\Psi(\vec{x}, \alpha, t), \Psi^+(\vec{x}', \alpha', t)]_{\pm} = \delta(\vec{x} - \vec{x}') \delta_{\alpha\alpha'}$$
$$N = \int_{\vec{x}, \alpha} \Psi^+(\vec{x}, \alpha, t) \Psi(\vec{x}, \alpha, t) = \sum_n a_n^+ a_n$$

Now we want introduce in the QED the essential role played by e.m. interactions in determining the structure and the properties of the innumerable systems of Condensed Matter



$$\vec{p}_i \rightarrow \vec{p}_i + e_i \vec{A}(\vec{x}_i, t)$$



$$H_{\text{matt}} = \sum_{j=1}^N H_0(\vec{x}_j, \vec{p}_j, \alpha_j) + H_{\text{rad}}^1 + H_{\text{rad}}^2 + V(\vec{x}_1, \alpha_1, \dots, \vec{x}_N, \alpha_N)$$

Coherence QED

$$H_{\text{matt}} = \sum_{j=1}^N H_0(\vec{x}_j, \vec{p}_j, \alpha_j) + H_{\text{rad}}^1 + H_{\text{rad}}^2 + H_{\text{SR}}$$

These two radiative terms make the standard QED a new QED

$$H_{\text{rad}}^1 = e \int_{\vec{x}, \alpha} \vec{A}(\vec{x}, t) \Psi^\dagger(\vec{x}, \alpha, t) \vec{J}(\alpha) \Psi(\vec{x}, \alpha, t)$$

$$H_{\text{rad}}^2 = e^2 \lambda \int_{\vec{x}, \alpha} \vec{A}^2(\vec{x}, t) \Psi^\dagger(\vec{x}, \alpha, t) \Psi(\vec{x}, \alpha, t)$$

$$H_{\text{SR}} = V(\vec{x}_1, \alpha_1, \dots, \vec{x}_N, \alpha_N)$$

Z-electron example for understand the role of λ

$$H = \sum_{r=1, Z} \frac{(p_r - eA(x_r, t))^2}{2m_e}$$



$$H_0 - \frac{e}{m_e} \sum_{r=1, Z} p_r A(x_r, t) + \frac{e^2}{2m_e} \sum_{r=1, Z} A^2(x_r, t)$$



$$J = - \sum_{r=1, Z} \frac{p_r}{m_e}$$

$$\lambda = \frac{Z}{2m_e}$$

$$H = H_0 + eJA(x_r, t) + e^2 \lambda A^2(x_r, t)$$

Coherence QED

Lagrangian of E.M. Field

$$\vec{A}(\vec{x}, t) = \sum_{\vec{k}r} \frac{1}{\sqrt{2\omega_{\vec{k}}V}} \left[a_{\vec{k}r}(t) \vec{\epsilon}_{\vec{k}r} e^{-i\omega_{\vec{k}}t} e^{i\vec{k}\cdot\vec{x}} + a_{\vec{k}r}^*(t) \vec{\epsilon}_{\vec{k}r}^* e^{i\omega_{\vec{k}}t} e^{-i\vec{k}\cdot\vec{x}} \right]$$

$$L_{em} = \sum_{\vec{k}r} \left[\frac{i}{2} (a_{\vec{k}r}(t)^* \dot{a}_{\vec{k}r}(t) - \dot{a}_{\vec{k}r}(t)^* a_{\vec{k}r}(t)) + \frac{1}{2\omega_{\vec{k}}} \dot{a}_{\vec{k}r}(t)^* \dot{a}_{\vec{k}r}(t) \right]$$

$$L_{matt} = i \int_{\vec{x}, \alpha} \Psi^+(\vec{x}, \alpha, t) \Psi_t(\vec{x}, \alpha, t) - H_{matt}$$

$$L_{QED-Coherence} = L_{matt} + L_{e.m.}$$

COHERENCE QED IN COLD FUSION

The existence of a large number N suggest rescaling theory so as to make N appear explicitly in the path-integral. We thus define:

$$\Psi_0(\vec{x}, \alpha, t) = \frac{1}{\sqrt{N}} \Psi(\vec{x}, \alpha, t) \quad a_{\vec{kr}}^0 = \frac{1}{\sqrt{N}} a_{\vec{kr}} \quad \int_{\vec{x}, \alpha} \Psi_0^+(\vec{x}, \alpha, t) \Psi_0(\vec{x}, \alpha, t) = 1$$

After rescaling of the lagrangians we can derive the motion equation (by path integral approach) by means of the principle of extremal action

$$\left\{ \begin{array}{l} \Psi_0(\vec{x}, \alpha, t) = \varphi(\vec{x}, \alpha, t) + \frac{1}{\sqrt{N}} \eta(\vec{x}, \alpha, t) \\ a_{\vec{kr}}^0 = \alpha_{\vec{kr}} + \frac{1}{\sqrt{N}} \beta_{\vec{kr}}(t) \end{array} \right. \longrightarrow \delta \int_{t_i}^{t_f} (L_{\text{matt}} + L_{\text{em}}) = 0$$

Coherence equation

$$\left\{ \begin{array}{l} i \frac{\partial \varphi}{\partial t} = H_0 \varphi + e \sqrt{N} \vec{A}_0 \cdot \vec{J} \varphi \\ -\frac{1}{2\omega_{\vec{k}}} \ddot{\alpha}_{\vec{kr}} + i \dot{\alpha}_{\vec{kr}} - \frac{e^2}{\omega_{\vec{k}}} \left(\frac{N}{V} \right) \lambda \alpha_{\vec{kr}} = e \frac{1}{\sqrt{2\omega_{\vec{k}}}} \left(\frac{N}{V} \right)^{1/2} \vec{\varepsilon}_{\vec{kr}}^* e^{i\omega_{\vec{k}} t} \int_{\vec{x}, \alpha} e^{-i\vec{k}\vec{x}} \varphi^* \vec{J} \varphi \end{array} \right.$$

QED-Coherence of TWO-LEVEL SYSTEMS

In this case we have:

$$\{\varphi_n(\vec{x}, \alpha)\} = \{\varphi_1(\vec{x}, \alpha), \varphi_2(\vec{x}, \alpha)\} \quad \begin{cases} H\varphi_1 = E_1\varphi_1 \\ H\varphi_2 = E_2\varphi_2 \end{cases} \quad E_1 - E_2 = \omega_0$$

Put $J_i = \delta_{il} J \sigma_l$ and using the interaction representation $\varphi_i(\vec{x}, t) = e^{-iE_i t} \chi_i(\vec{x}, t)$

we obtain the following *coherence equation*

$$i \frac{\partial \chi_2(\vec{x}, t)}{\partial t} = -\frac{\vec{\nabla}^2}{2m} \chi_2(\vec{x}, t) + eJ \sqrt{\frac{N}{V}} \frac{1}{\sqrt{2\omega_0}} \sum_{|\vec{k}|=\omega_0, r} \vec{\varepsilon}_{\vec{k}r}^* \alpha_{\vec{k}r}^* e^{-i\vec{k}\vec{x}} \chi_1(\vec{x}, t)$$

$$i \frac{\partial \chi_1(\vec{x}, t)}{\partial t} = -\frac{\vec{\nabla}^2}{2m} \chi_1(\vec{x}, t) + eJ \sqrt{\frac{N}{V}} \frac{1}{\sqrt{2\omega_0}} \sum_{|\vec{k}|=\omega_0, r} \vec{\varepsilon}_{\vec{k}r}^* \alpha_{\vec{k}r}^* e^{-i\vec{k}\vec{x}} \chi_2(\vec{x}, t)$$

$$-\frac{1}{2\omega_0} \ddot{\alpha}_{\vec{k}r} + i\dot{\alpha}_{\vec{k}r} - \frac{e^2}{\omega_0} \left(\frac{N}{V}\right) \lambda \alpha_{\vec{k}r} = eJ \frac{1}{\sqrt{2\omega_0}} \left(\frac{N}{V}\right)^{1/2} \vec{\varepsilon}_{\vec{k}r}^* \int_{\vec{x}} e^{-i\vec{k}\vec{x}} \chi_2^* \chi_1$$

$$\alpha_{\vec{k}r} = 0 \quad \text{for} \quad |\vec{k}| \neq \omega_0$$

Coherence domains

In the last equation, we can observe that the resonating modes whose vacuum frequencies $\omega_k = \omega_0$ are strongly coupled with the matter field χ . It produce the remarkable result that the e.m. and matter field exhibit a space-structure comprising an array of COHERENCE DOMAINS (CD)

$$CD \text{ size} = 2\pi/\omega_0$$

Now if we work within a CD, we can write $\chi_i(\vec{x}, t) \approx \frac{1}{\sqrt{V_{CD}}} \chi_i(t)$ so we have the following coherence equation for a CD:

$$i\dot{\chi}_2(t) = eJ \sqrt{\frac{N}{V}} \frac{1}{\sqrt{2\omega_0}} \sum_r \int d\Omega_{\vec{k}} \varepsilon_{\vec{k}r,1}^* \alpha_{\vec{k}r}^* \chi_1(t)$$

$$i\dot{\chi}_1(t) = eJ \sqrt{\frac{N}{V}} \frac{1}{\sqrt{2\omega_0}} \sum_r \int d\Omega_{\vec{k}} \varepsilon_{\vec{k}r,1} \alpha_{\vec{k}r} \chi_2(t)$$

$$-\frac{1}{2\omega_0} \ddot{\alpha}_{\vec{k}r} + i\dot{\alpha}_{\vec{k}r} - \frac{e^2}{\omega_0} \left(\frac{N}{V}\right) \lambda \alpha_{\vec{k}r} = eJ \frac{1}{\sqrt{2\omega_0}} \left(\frac{N}{V}\right)^{1/2} \vec{\varepsilon}_{\vec{k}r,1}^* \chi_2^*(t) \chi(t)_1$$

$$|\chi_1|^2 + |\chi_2|^2 = 1$$

Coherence domains

Putting

$$A = \sum_r \left(\frac{3}{8\pi} \right)^{1/2} \int d\Omega_{\vec{k}} \varepsilon_{\vec{k}r,1} \alpha_{\vec{k}r} \quad g = eJ \left(\frac{8\pi}{3} \right)^{1/2} \left(\frac{N}{2V\omega_0^3} \right)^{1/2} \quad \mu = \frac{e^2 \lambda}{\omega_0^2} \left(\frac{N}{V} \right)$$

We can write the coherence equation within CD in this way

$$\dot{\chi}_2 = -igA^* \chi_1 \quad \dot{\chi}_1 = -igA \chi_2 \quad \frac{i}{2} \ddot{A} + \dot{A} + i\mu A = -ig\chi_2^* \chi_1$$

If we derive the e.m. equation we will have

$$\frac{i}{2} \ddot{A} + \ddot{A} + i\mu \dot{A} + gA = 0$$

$$\frac{p^3}{2} - p^2 - \mu p + g^2 = 0 \quad g^2 > g_c^2 \quad g_c^2 = \frac{8}{27} + \frac{2}{3} \mu + \left[\frac{4}{9} + \frac{2}{3} \mu^2 \right]^{3/2}$$

Finally we have these coserved quantities

$$Q = A^* A + \frac{i}{2} (A^* \dot{A} - \dot{A}^* A) + \chi_1^* \chi_1 \quad H = Q + \frac{1}{2} \dot{A}^* \dot{A} + \mu A^* A + g (A^* \chi_2^* \chi_1 + A \chi_1^* \chi_2)$$

$$1 = \chi_1^* \chi_1 + \chi_2^* \chi_2$$

Coherence QED of a plasma

$$H_0 = \sum_i \frac{\vec{p}_i^2}{2m} + \frac{1}{2} m \omega_p \xi_i^2$$

$$H_{rad}^1 = e \int \Psi^\dagger(\vec{x}, \xi, t) \vec{J} \Psi(\vec{x}, \xi, t) \vec{A}(\vec{x}, t)$$

$$J = Q(p/m)$$

If we resolve the CE for a plasma we obtain with $\mu=0$

$$g^2 = \left(\frac{2\pi}{3}\right) \quad \longrightarrow \quad g_c^2 = \frac{16}{27} < g^2 = \left(\frac{2\pi}{3}\right)$$

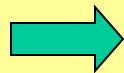
But this solution don't satisfy the conserved quantities

$$Q = \sum_k \left(A_k^* A_k + \frac{i}{2} (A_k^* \dot{A}_k - \dot{A}_k^* A_k) + \chi_k^* \chi_k \right)$$

$$H = \frac{E}{N\omega_p} = Q + \sum_k \left[\frac{1}{2} \dot{A}_k^* \dot{A}_k - ig (A_k^* \chi_k - A_k \chi_k^*) \right]$$

$$\chi_k \prec e^{i\psi}$$

$$A_k \prec e^{i\phi}$$



$$1 - \dot{\phi} + \frac{g^2}{\dot{\phi}^2} = 0$$

No Solution!

Coherence QED of a plasma

We must introduce another interaction term

$$H_0^R = \sum_i \frac{\vec{p}_i^2}{2m} + \frac{1}{2} m \omega_p \xi_i^2 + \omega_R \left(\sum_k \frac{a_k^+ a_k}{\alpha_{\max}^2} \right)$$

$$\alpha_{\max} = \sqrt{m \omega_p} \left(\frac{V}{N} \right)^{1/3} \quad \text{and} \quad \omega_R = 1.88 \omega_p$$

The effect of this term is to strongly suppress coherent oscillation amplitudes when

$$\alpha^2 \geq \alpha_{\max}^2$$

In this way, we obtain as solution a e.m. field with a new frequency

$$\omega_r = \omega_R (1 - \dot{\phi}) \quad \longrightarrow \quad \dot{\phi} = 1 + \sqrt{1 - 2g \frac{\chi}{A}}$$

Cold fusion phenomenon

Why the D_2 on the PD surface became D_2^+ ?

Now we propose an explanation about the remarkable property of Pd to ionize D_2 or H_2 gas (need 30 eV!)

Due to the plasma in the lattice (10 electrons in the d-band) we have an e.m. field

$$\nabla^2 \vec{A} = 0$$

If we work in semi-space $z > 0$

$$\vec{A}(x, y, z) = \vec{A}(x, y, 0)e^{-\eta z} \quad \eta = \sqrt{\omega_p^2 - \omega_r^2} > 0$$

If we suppose that the system D_2 -molecule can be reduced at a two-dimensional system ground and ionized we have (\vec{D} is the dipole electric operator):

$$|\chi\rangle = c_1|0\rangle + c_2|1\rangle \quad \text{and this interaction term} \quad H_{\text{int}} = -\vec{E} \cdot \vec{D}$$

Cold fusion phenomenon

So put

$$\vec{E} \approx \vec{u} \sin(\omega_r t) \sqrt{\frac{2N\omega_r^2}{V\omega_p}} \sqrt{\frac{2\pi}{3}} A \quad \omega^* = \sqrt{\frac{2N\omega_r^2}{V\omega_p}} \sqrt{\frac{2\pi}{3}} A e |\langle i | H_{\text{int}} | 0 \rangle| \quad |\langle i | H_{\text{int}} | 0 \rangle| = 1 \text{ \AA}$$

$$i \frac{\partial}{\partial t} |\chi\rangle = \left[\frac{\omega_i}{2} \sigma_3 + \frac{\omega^*}{2} (e^{-i\omega_r t} \sigma_+ + e^{-i\omega_r t} \sigma_-) \right] |\chi\rangle$$

$$\omega^* \approx 10 \text{ eV} \quad \omega_i \approx 30 \text{ eV} \quad \omega_r \approx 0.2 \omega_p \quad \frac{N}{V} = 6.5 \cdot 10^{22} \text{ cm}^{-3} \quad \omega_p \approx 10 \text{ eV}$$

$$P_i(t) = \frac{\omega^{*2}}{(\omega_i - \omega_r)^2 + \omega^{*2}} \sin^2 \sqrt{(\omega_i - \omega_r)^2 + \omega^{*2}} \frac{t}{2}$$

The average probability that a D₂-molecule be ionised on the surface of Pd is of the order of 5%!

Tunneling Effect Probability within lattice

$$|P|_{\text{int}}^2 = \exp\left(-2 \int_0^\alpha K(r)_{\text{int}} dr\right) \quad K(r)_{\text{int}} = \sqrt{2\mu[E - V(r)]/\hbar^2}$$

where α is about 0.10 \AA . In this case:

$$V(r)_{\text{int}} = k \frac{q^2}{r} \cdot \left[V(r)_M - J \frac{\xi \hbar \omega_r R}{r} \right] \quad \longrightarrow \quad \text{is the contribute due to QED-Coherence plasma}$$

R is the nuclear radius, ξ a parameter varying between 0.015 and 0.025 which depends on the structural characteristics of the lattice, T the absolute temperature at which the metal is placed experimentally.

In this case, the Morse potential will be:

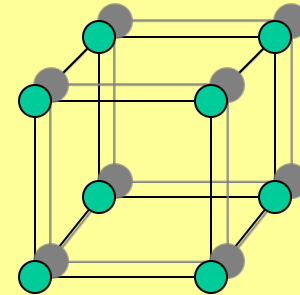
$$V(r)_M = \left(J / \zeta \right) \exp\left(-2 \varphi\left(r - r_0 \right) \right) - 2 \exp\left(-\varphi\left(r - r_0 \right) \right)$$

φ and ζ depend on the dynamic conditions of the system, while r_0 is the point of classic inversion.

Cold fusion phenomenon

J = impurities concentration

$$\lambda_f = \lambda \frac{4\pi\rho\hbar}{\mu_d} \left\langle \frac{1}{p} \right\rangle = \text{rate fusion for energy}$$



$$D_L = \iiint_{\omega} J \frac{\rho l^2 v b^2}{\alpha 2hR} \exp\left(-\frac{U_0}{kT}\right) d\omega = \text{3D lattice deformation}$$

$$\Gamma \approx \frac{\exp\left(-2 \int_0^{\alpha} K(r)_{\text{int}} dr\right)}{\lambda \cdot \frac{4\pi\rho\hbar}{\mu_d} \cdot \left\langle \frac{1}{p} \right\rangle} \cdot D_L$$

Cold fusion phenomenon

Tunneling Effect probability + 3D lattice interaction

Palladium $J \approx 0.75\%$ $T(K) \approx 100-300K$ $\alpha \approx 10\text{\AA}$ $\lambda = 10^{-3} \text{ eV / min}$ $M_{Pd} / (\mu g)$

T \approx	100 K	140 K	180 K	220 K	260 K	300 K
E (eV)	$\Gamma \approx$	$\Gamma \approx$	$\Gamma \approx$	$\Gamma \approx$	$\Gamma \approx$	$\Gamma \approx$
150	10^{-77}	10^{-67}	10^{-71}	10^{-63}	10^{-60}	10^{-59}
160	10^{-73}	10^{-65}	10^{-70}	10^{-61}	10^{-56}	10^{-55}
170	10^{-71}	10^{-62}	10^{-69}	10^{-58}	10^{-55}	10^{-52}
180	10^{-69}	10^{-60}	10^{-67}	10^{-54}	10^{-53}	10^{-50}
190	10^{-68}	10^{-54}	10^{-64}	10^{-53}	10^{-52}	10^{-45}
200	10^{-64}	10^{-53}	10^{-63}	10^{-51}	10^{-50}	10^{-43}
210	10^{-63}	10^{-54}	10^{-64}	10^{-50}	10^{-48}	10^{-40}
220	10^{-61}	10^{-52}	10^{-63}	10^{-49}	10^{-45}	10^{-35}
230	10^{-60}	10^{-50}	10^{-62}	10^{-47}	10^{-43}	10^{-29}
240	10^{-59}	10^{-45}	10^{-60}	10^{-44}	10^{-39}	10^{-25}
250	10^{-58}	10^{-41}	10^{-59}	10^{-40}	10^{-36}	10^{-21}

Morse potential

Conclusion

- *Introducing in the QED the essential role played by e.m. interactions in the Condensed Matter we can predict new phenomenon*
- *The Coherence Domains are the new structures of Condensed Matter*
- *By means of this theory we can explain the ionisation of H_2 or D_2 molecule on the Pd surface*
- *Integrating this theory with the microcrack formation we are able to explain the cold fusion phenomenon*